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# COMPUTATIONAL DYNAMICS OF ELASTOMERIC-BASED ISOLATION SYSTEMS FOR RIGID STRUCTURES

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Abstract: The paper analyses the dynamic behaviour of a rigid structure elastic insulated, subjected to various actions of vibration or shock type. We have considered for analysis the classical model of a rigid structure with six degree of freedom in two cases: with one vertical symmetry plane and with two vertical symmetry planes. For the chosen model was performed modal analysis for a set of rigidities containing six different values. The study aims to highlight the correlation between the isolation characteristics and natural pulsations of the structure in terms of eigenvectors and eigenvalues in the indicated analysis conditions.

Keywords: computational mechanics, eigenvalue, comparative analysis, eigenvector, dynamic model, natural pulsations

#### 1. INTRODUCTION

Dynamic behavior of structural systems is described by mathematical equations that take different forms from a specific case to another. The representation in calculations of the solid bodies (rigid or deformable), of the linkages and of the whole system has at base the concept of dynamic system and dynamic model.

The dynamic system is an abstraction of the physical and mechanical characteristics of the structural system whose mechanical condition changes during time.

Any dynamical system has characterized by some specific qualitative properties (inertial, dissipative, and elastic) represented by the values of measurable parameters (mass, moments of inertia, the damping coefficient, the rigidity/ flexibility coefficient).

The dynamic model is essentially an idealized form, simplified or schematized of a dynamic system in order to reduce the numerical analysis operations without that the real processes (qualitative and quantitative) to be significantly modified.

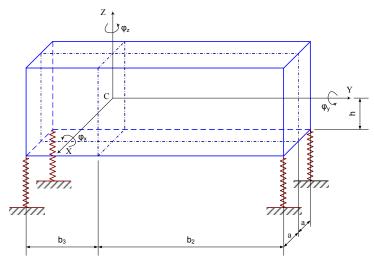
The dynamic response is the instantaneous state of a dynamic system over which have been applied external dynamic actions, real and variable during time. The dynamic response can be expressed through fundamental kinematic parameters (displacement, velocity, acceleration) or through derived parameters (energy, sectional strains, stresses, deformations, generalized forces).

#### 2. THEORETICAL APPROACHES

It is proposed the complex dynamic model of a rigide structure with six degrees of freedom which consisting of three translational linear coordinates x, y, z and three rotational angular coordinates  $\varphi_x$ ,  $\varphi_y$ ,  $\varphi_z$ . For the proposed model, we will study the behavior of the structure under the vibration action and in the presence of elastic elements.

Thus, the motion equations of the rigid with elastic linkages are written:

$$\underline{\underline{A}}_{\underline{\underline{q}}}^{\underline{q}} + \underline{\underline{C}}_{\underline{\underline{q}}}^{\underline{q}} = \underline{\underline{0}}$$
 (1)



**Figure 1:** Model of the rigid with six degrees of freedom, elastic supported in four points on inferior base, with a longitudinal vertical plane of symmetry

In an analytical form, the system is:

$$\begin{cases} m\ddot{x} + x \sum k_{ix} + \varphi_{y} \sum k_{ix} z_{i} - \varphi_{z} \sum k_{ix} y_{i} = 0 \\ m\ddot{y} + y \sum k_{iy} - \varphi_{x} \sum k_{iy} z_{i} + \varphi_{z} \sum k_{iy} x_{i} = 0 \\ m\ddot{z} + z \sum k_{iz} + \varphi_{x} \sum k_{iz} y_{i} - \varphi_{y} \sum k_{iz} x_{i} = 0 \end{cases}$$

$$J_{x}\ddot{\varphi}_{x} - y \sum k_{iy} z_{i} + z \sum k_{iz} y_{i} + \varphi_{x} \sum (k_{iy} z_{i}^{2} + k_{iz} y_{i}^{2}) - \varphi_{y} \sum k_{iz} x_{i} y_{i} - \varphi_{z} \sum k_{iy} z_{i} x_{i} = 0$$

$$J_{y}\ddot{\varphi}_{y} + x \sum k_{ix} z_{i} - z \sum k_{iz} x_{i} - \varphi_{x} \sum k_{iz} x_{i} y_{i} + \varphi_{y} \sum (k_{iz} x_{i}^{2} + k_{ix} z_{i}^{2}) - \varphi_{z} \sum k_{ix} y_{i} z_{i} = 0$$

$$J_{z}\ddot{\varphi}_{z} - x \sum k_{ix} y_{i} + y \sum k_{iy} x_{i} - \varphi_{x} \sum k_{iy} z_{i} x_{i} - \varphi_{y} \sum k_{ix} y_{i} z_{i} + \varphi_{z} \sum (k_{ix} y_{i}^{2} + k_{iy} x_{i}^{2}) = 0$$

$$(2)$$

The system (2) is difficult to solve analytically or using the matriceal formalism because it requires a large amount of calculation, and the sixth degree polynomial equation of the natural pulsations involves difficulties in solving and analysis. The solution is the automatic numerical calculus of the differential motion equations system of second degree, resulting a system with 12 differential equations of first degree, which can be integrated without difficulty. On the other hand, at the use of numerical analysis appears as a disadvantage the highlighting of the influence of the dynamic system physical characteristics. Thus, the analysis is done by repeated tests, using different sets of values for the input data. To this end, both for the elimination of coupling movements and to analytical solve of the dynamic system model, may impose certain sized and structure requirements for the system, leading to a decoupling of the equation system into subsystems simple and easier to integrate.

As discussed above, we consider the case of the rigid structure, elastic supported in four points on inferior base, with a longitudinal vertical plane of symmetry yCz as in Figure 1. In this situation a few simplifying assumptions are valid:

- the dimensions of the analyzed rigid structure are symmetrical in relation to the considered plane
- the elastic linkages are identical, have symmetrical positions and are located in the same horizontal plane Due to the mentioned symmetries, a part of coupling terms from stiffness matrix are canceled, and we have:

$$\sum k_{iy} x_i = 0$$

$$\sum k_{iz} x_i = 0$$

$$\sum k_{iz} x_i y_i = 0$$

$$\sum k_{iy} z_i x_i = 0$$
(3)

Through the disappearance of the coupling terms, the system decouples into two subsystems described by coordinates  $(y, z, \varphi_x)$  and  $(x, \varphi_v, \varphi_z)$ .

For the two decoupled subsystems can be written the equations of the free vibrations. Thus, we have for the subsystem  $(y, z, \varphi_x)$ :

$$\begin{cases} m\ddot{y} + 4k_{y}y + 4hk_{y}\varphi_{x} = 0\\ m\ddot{z} + 4k_{z}z + 2k_{z}(b_{3} - b_{2})\varphi_{x} = 0\\ J_{x}\ddot{\varphi}_{x} + 4hk_{y}y + 2k_{z}(b_{3} - b_{2})z + 2[k_{z}(b_{2}^{2} + b_{3}^{2}) + 2h^{2}k_{y}]\varphi_{x} = 0 \end{cases}$$

$$(4)$$

and for the subsystem  $(x, \varphi_y, \varphi_z)$ :

$$\begin{cases} m\ddot{x} + 4k_{x}x - 4hk_{x}\varphi_{y} - 2k_{x}(b_{3} - b_{2})\varphi_{z} = 0\\ J_{y}\ddot{\varphi}_{y} - 4hk_{x}x + 4(h^{2}k_{x} + a^{2}k_{z})\varphi_{y} + 2hk_{x}(b_{3} - b_{2})\varphi_{z} = 0\\ J_{z}\ddot{\varphi}_{z} - 2k_{x}(b_{3} - b_{2})x + 2hk_{x}(b_{3} - b_{2})\varphi_{y} + 2[2a^{2}k_{y} + k_{x}(b_{2}^{2} + b_{3}^{2})]\varphi_{z} = 0 \end{cases}$$

$$(5)$$

Further it is proposed an analyze of the vibrations of the two subsystems characterized each by three dynamic coordinates (degrees of freedom) coupled.

For each of the two subsystems, with elastic linkages and three degrees of freedom, the vector of the generalized coordinates is:

$$q = [q_1, q_2, q_3]^T \tag{6}$$

Using the classical mathematical apparatus, were writing the quadratic forms of the system energies and then the case II Lagrange equations were used for the obtaining of the motion equations, written matriceal in the form (1). The solution for the system (1) has been sought as:

(1). The solution for the system (1) has been sought as:
$$\frac{q}{=} = \underbrace{a \sin pt}_{=} = \begin{cases} a_1 \\ a_2 \\ a_3 \end{cases} \sin pt$$
where  $\underbrace{a}_{=} = \begin{cases} a_1 \\ a_2 \\ a_3 \end{cases}$  is the vector of the motion amplitudes.

Taking into account the proposed form of the system solution, equation (1) becomes:

$$(\underline{C} - p^2 \underline{A})\underline{a} = \underline{0} \tag{8}$$

Equation (8) has nonzero solutions only if the determinant of the matrix is zero. This equation has third degree with p<sub>2</sub> as variable and represent the equation of the natural pulsations of the dynamic system with three degrees of freedom. Through the resolving analytically or numerically of the equation (8), we obtain the three natural pulsations of the system  $p_1$ ,  $p_2$ ,  $p_3$ .

#### 3. CASE STUDY

The case study was made for two types of symmetry of the proposed structure, namely:

 considering the structure having a longitudinal vertical plane of symmetry, case for which as numerical values were proposed:

a=7.5 m  $b_3 = 12 \text{ m}$  $b_2=8 \text{ m}$ h=7 m  $J_x = 42 \times 10^6 \text{ kgm}^2$  $J_v = 25 \times 10^6 \text{ kgm}^2$  $J_z = 17.5 \times 10^6 \text{ kgm}^2$ 

 considering the structure with two vertical planes of symmetry, one longitudinal and one transversal, case for which as numerical values were proposed:

a=7.5 m  $b_3 = 10 \text{ m}$  $b_2 = 10 \text{ m}$ h=7 m  $J_x = 35 \times 10^6 \text{ kgm}^2$  $J_{v} = 25 \times 10^{6} \text{ kgm}^{2}$  $J_z = 17.5 \times 10^6 \text{ kgm}^2$  For both sets of values, the mass of the analyzed structure was considered with the value m = 3x106 kg. Also, in both cases were proposed for study six sets of values of stiffness coefficients as follows in Table 1.

**Table 1:** The proposed sets of values for the stiffness coefficients

	Var 1	Var 2	Var 3	Var 4	Var 5	Var 6
k <sub>x</sub> [N/m]	$2x10^{6}$	$4x10^{6}$	$8x10^{6}$	$16x10^{6}$	$32x10^{6}$	$64x10^6$
$k_v[N/m]$	$2x10^{6}$	$4x10^{6}$	$8x10^{6}$	$16x10^6$	$32x10^{6}$	$64x10^6$
$k_z[N/m]$	$8x10^{6}$	$16x10^{6}$	$32x10^{6}$	$64x10^{6}$	$128x10^6$	$256x10^6$

After completing the mathematical calculus, we obtain values for the parameters initially proposed - the eigenvalues and eigenvectors of the analyzed system. These values are summarized in Tables 2, 3, 4, 5.

**Table 2:** The system parameters assessment for the subsystem I in the case with one vertical plane of symmetry

Table	Table 2: The system parameters assessment for the subsystem 1 in the case with one vertical plane of symmetry									
	Case with a longitudinal vertical plane of symmetry - subsystem I									
	eigen	nat puls	freq	eigen	eigen	eigen	eigen	eigen	eigen	
	val	p	f	vect 1	vect 2	vect 3	vect 1	vect 2	vect 3	
	$p^2$			$\mu_1$	$\mu_2$	$\mu_3$	norm	norm	norm	
							$\mu_{ln}$	$\mu_{2n}$	$\mu_{3\mathrm{n}}$	
	89.3	9.4484	1.5038	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V1	2.4	1.5376	0.2447	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	10.3	3.2044	0.5100	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	
	178.5	13.3621	2.1266	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V2	4.7	2.1745	0.3461	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	20.5	4.5317	0.7212	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	
	357.1	18.8968	3.0075	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V3	9.5	3.0752	0.4894	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	41.1	6.4088	1.0200	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	
	714.2	26.7241	4.2533	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V4	18.9	4.3490	0.6922	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	82.1	9.0634	1.4425	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	
	1428.4	37.7936	6.0150	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V5	37.8	6.1504	0.9789	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	164.3	12.8176	2.0400	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	
	2856.7	53.4483	8.5066	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V6	75.7	8.6980	1.3843	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	328.6	18.1268	2.8850	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	

Table 3: The system parameters assessment for the subsystem II in the case with one vertical plane of symmetry

	Case with a longitudinal vertical plane of symmetry - subsystem II										
	eigen	nat puls	freq	eigen	eigen	eigen	eigen	eigen	eigen		
	val	p	f	vect 1	vect 2	vect 3	vect 1	vect 2	vect 3		
	$p^2$			$\mu_1$	$\mu_2$	$\mu_3$	norm	norm	norm		
							$\mu_{1n}$	$\mu_{2n}$	$\mu_{3n}$		
	2.1	1.4601	0.2324	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000		
V1	90.0	9.4861	1.5098	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657		
	71.5	8.4550	1.3457	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837		
	4.3	2.0649	0.3286	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000		
V2	180.0	13.4153	2.1351	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657		
	143.0	11.9571	1.9030	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837		
	8.5	2.9202	0.4648	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000		
V3	359.9	18.9721	3.0195	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657		
	285.9	16.9100	2.6913	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837		
	17.1	4.1298	0.6573	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000		
V4	719.9	26.8306	4.2702	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657		
	571.9	23.9143	3.8061	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837		
	34.1	5.8405	0.9295	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000		

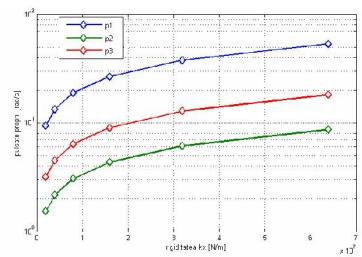
	1439.8	37.9442	6.0390	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657
	1143.8	33.8199	5.3826	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837
	68.2	8.2597	1.3146	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000
V6	2879.5	53.6612	8.5405	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657
	2287.6	47.8286	7.6122	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837

Table 4: The system parameters assessment for the subsystem I in the case with 2 vertical planes of symmetry

	Case with two vertical planes of symmetry (longitudinal and vertical) - subsystem I								
	eigenval	nat puls	freq	eigenvect 1	eigenvect 2	eigenvect 3			
	$p^2$	p	f	$\mu_1$	$\mu_2$	$\mu_3$			
	2.4	1.5391	0.2450	-0.9999	-0.1830	0			
V1	102.9	10.1453	1.6147	0	0	1.0000			
	10.7	3.2660	0.5198	0.0160	-0.9831	0			
	4.7	2.1766	0.3464	-0.9999	-0.1830	0			
V2	205.9	14.3476	2.2835	0	0	1.0000			
	21.3	4.6188	0.7351	0.0160	-0.9831	0			
	9.5	3.0782	0.4899	-0.9999	-0.1830	0			
V3	411.7	20.2905	3.2293	0	0	1.0000			
	42.7	6.5320	1.0396	0.0160	-0.9831	0			
	19.0	4.3532	0.6928	-0.9999	-0.1830	0			
V4	823.4	28.6952	4.5670	0	0	1.0000			
	85.3	9.2376	1.4702	0.0160	-0.9831	0			
	37.9	6.1563	0.9798	-0.9999	-0.1830	0			
V5	1646.8	40.5811	6.4587	0	0	1.0000			
	170.7	13.0639	2.0792	0.0160	-0.9831	0			
	75.8	8.7064	1.3857	-0.9999	-0.1830	0			
V6	3293.6	57.3903	9.1340	0	0	1.0000			
	341.3	18.4752	2.9404	0.0160	-0.9831	0			

**Table 5:** The system parameters assessment for the subsystem II in the case with 2 vertical planes of symmetry

	Case wi	ith two vertical	planes of symn	netry (longitudinal a	and vertical) - subsyst	em II
	eigenval	nat puls	freq	eigenvect 1	eigenvect 2	eigenvect 3
	$p^2$	p	f	$\mu_1$	$\mu_2$	$\mu_3$
	2.2	1.4757	0.2349	-0.9997	0.2133	0
V1	88.2	9.3898	1.4944	-0.0262	-0.9770	0
	71.4	8.4515	1.3451	0	0	1.0000
	4.4	2.0869	0.3321	-0.9997	0.2133	0
V2	176.3	13.2792	2.1135	-0.0262	-0.9770	0
	142.9	11.9523	1.9023	0	0	1.0000
	8.7	2.9514	0.4697	-0.9997	0.2133	0
V3	352.7	18.7797	2.9889	-0.0262	-0.9770	0
	285.7	16.9031	2.6902	0	0	1.0000
	17.4	4.1739	0.6643	-0.9997	0.2133	0
V4	705.4	26.5585	4.2269	-0.0262	-0.9770	0
	571.4	23.9046	3.8045	0	0	1.0000
	34.8	5.9027	0.9394	-0.9997	0.2133	0
V5	1410.7	37.5593	5.9778	-0.0262	-0.9770	0
	1142.9	33.8062	5.3804	0	0	1.0000
	69.7	8.3477	1.3286	-0.9997	0.2133	0
V6	2821.4	53.1169	8.4538	-0.0262	-0.9770	0
	2285.7	47.8091	7.6091	0	0	1.0000



**Figure 2:** The dependence between the natural pulsations of the system and the rigidity in the horizontal direction

Based on the values obtained for the three natural pulsations of the system in the six cases proposed, it was realized a graphic (Figure 2) of the dependence between the natural pulsations and the rigidity in the horizontal direction, denoted by  $k_x$ .

#### 4. CONCLUSIONS

It mentions that the stiffness in the x direction is equal to that in the y direction and stiffness in the z direction is a linear combination of the two others. Therefore we obtained for the eigenvectors identical values, regardless of the values of the stiffness coefficient considered in calculus.

As an independent variable for the representation of the pulsation evolution was chosen the stiffness in the x direction, denoted by  $k_x$ .

The evolution of the pulsations corresponding to the eigenvalues follows the natural tendency imposed by the pairs of values considered for rigidities. The correlative analysis of each set of eigenvalues induce the following conclusion, namely that linear combinations between the stiffness in the horizontal plane and the one in the vertical plane require similar evolutions.

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